

Boundary Conditions

Lecture 13

Boundary Conditions

- Boundary conditions are the relations between the electromagnetic field vectors on two sides of the interface that separates the two media.
- These conditions are dictated by types of material by which media are made of.

Electrostatic Boundary Conditions

- **At any point on the boundary,**
 - the components of \mathbf{E}_1 and \mathbf{E}_2 tangential to the boundary are equal

$$E_{1t} = E_{2t}$$

- the components of \mathbf{D}_1 and \mathbf{D}_2 normal to the boundary are equal

$$D_{1n} = D_{2n}$$

Magnetic Boundary Conditions

- The normal component of \mathbf{B}_1 and \mathbf{B}_2 continuous across a interface:

$$B_{1n} = B_{2n}$$

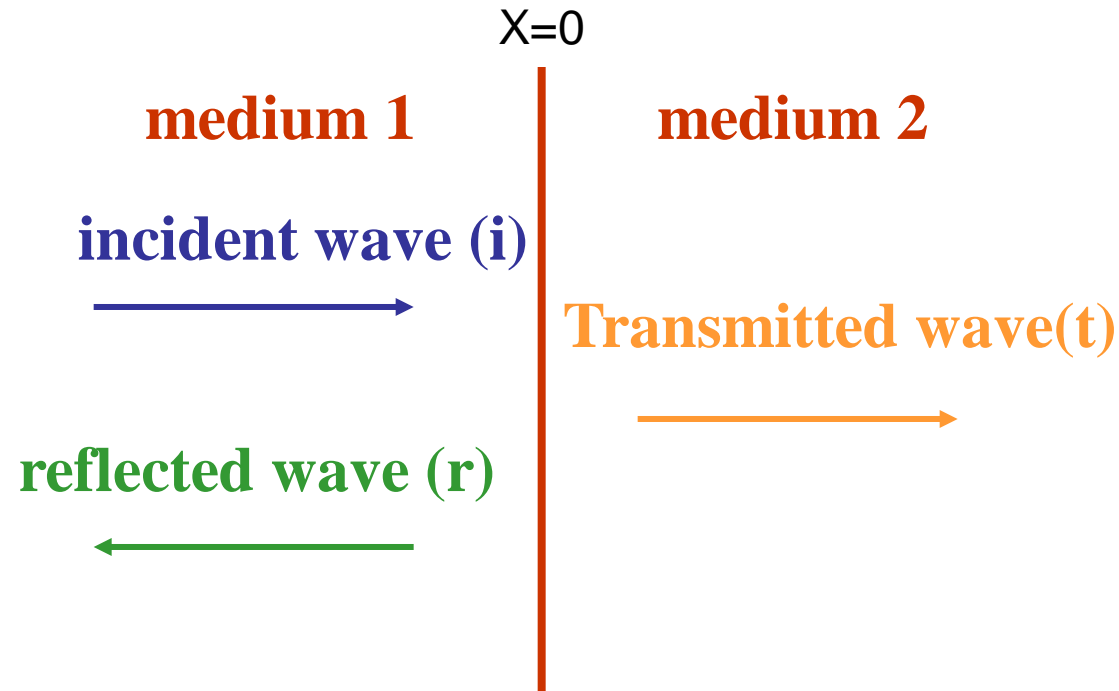
- The tangential component of a \mathbf{H}_1 and \mathbf{H}_2 to the boundary are equal

$$H_{1t} = H_{2t}$$

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Reflection and Transmission of Waves at Planar Interfaces

1. Reflection of uniform plane waves by perfect dielectric for normal incidence.



A traveling E wave approaches the interface $x=0$ from region 1, $x < 0$.

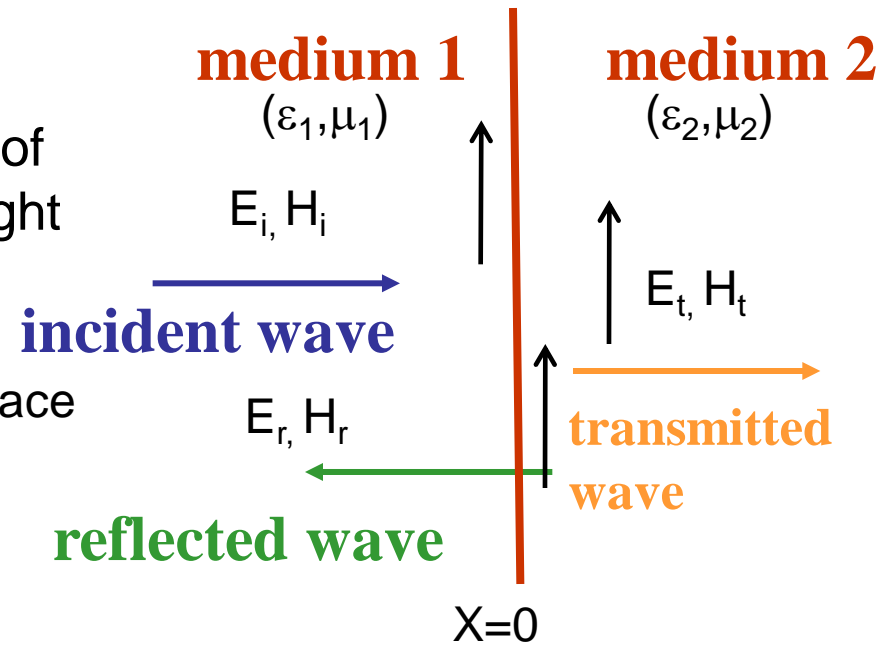
E_i and E_r are at $x = -0$, while E_t is at $x = +0$

Consider the case of plane em wave travelling in x direction and polarized in y direction and is incident normally on the interface from left.

It give rise to a reflected wave which travels back to the left in medium 1 and transmitted wave which continues on the right in medium 2.

Let (ϵ_1, μ_1) and (ϵ_2, μ_2) are the dielectric permittivities and magnetic permeabilities of two medium 1 and medium 2 at left and right of the plane $x=0$.

In this case E and H are tangential to the interface



The intrinsic impedance for medium 1

$$Z_1 = \sqrt{\frac{\mu_1}{\epsilon_1}}$$

The intrinsic impedance for medium 2

$$Z_2 = \sqrt{\frac{\mu_2}{\epsilon_2}}$$

The electric and magnetic fields associated with incident, reflected and transmitted em waves are related as

$$E_i = Z_1 H_i \quad \dots 1$$

$$E_r = -Z_1 H_r \quad \dots 2$$

$$E_t = Z_2 H_t \quad \dots 3$$

According to the boundary conditions at the interface for the continuity of tangential components of E and H, the electric and magnetic fields on each side of the boundary are same.

$$E_i + E_r = E_t \quad \dots 4$$

$$H_i + H_r = H_t \quad \dots 5$$

from eq 1, 2 and 3

$$H_i = \frac{E_i}{Z_1}, \quad H_r = -\frac{E_r}{Z_1}, \quad H_t = \frac{E_t}{Z_2}$$

putting the value of H_i, H_r, H_t in eq 5

$$\frac{E_i}{Z_1} - \frac{E_r}{Z_1} = \frac{E_t}{Z_2} \quad \dots 6$$

Putting the value of E_t from eq 4 in eq 6

$$\frac{1}{Z_1} (E_i - E_r) = \frac{(E_i + E_r)}{Z_2}$$

$$E_i(Z_2 - Z_1) = E_r(Z_2 + Z_1)$$

$$\frac{E_r}{E_i} = \frac{(Z_2 - Z_1)}{(Z_2 + Z_1)} = R_E = \text{Reflection coefficient}$$

for electric field intensity7

Transmission coefficient for electric field intensity

$$\begin{aligned}T_E &= \frac{E_t}{E_i} = \frac{E_i + E_r}{E_i} \\ &= 1 + \frac{E_r}{E_i} = 1 + \frac{(Z_2 - Z_1)}{(Z_2 + Z_1)} \\ T_E &= \frac{2Z_2}{(Z_2 + Z_1)} \quad \dots 8\end{aligned}$$

Similarly the reflection and transmission coefficient for magnetic field intensity H is obtained

Reflection coefficient for magnetic field intensity

$$R_H = \frac{H_r}{H_i} = -\frac{E_r}{E_i} = \frac{(Z_1 - Z_2)}{(Z_1 + Z_2)} \quad \dots 9$$

Transmission coefficient for magnetic field intensity

$$T_H = \frac{H_t}{H_i} = \frac{E_t / Z_2}{E_i / Z_1} = \frac{Z_1 E_t}{Z_2 E_i} = \frac{Z_1}{Z_2} T_E$$

$$T_H = \frac{2Z_1 Z_2}{Z_2 (Z_2 + Z_1)} = \frac{2Z_1}{(Z_2 + Z_1)} \quad \dots 10$$

Eq (7),(8),(9) and (10) are valid for a boundary between any two materials

In general equations for E can be written as

$$\overline{E}_i(x, t) = E_{i0} e^{-\gamma_1 x} e^{j\omega t} \hat{a}_x$$

$$\overline{E}_r(x, t) = E_{r0} e^{-\gamma_1 x} e^{j\omega t} \hat{a}_x$$

$$\overline{E}_t(x, t) = E_{t0} e^{-\gamma_2 x} e^{j\omega t} \hat{a}_x$$

Here E_{i0} , E_{r0} and E_{t0} are the electric field intensities at $x=0$ and γ_1, γ_2 are the propagation constants in 1st and 2nd medium

The equations for H can be written as

$$\overline{H}_i(x, t) = H_{i0} e^{-\gamma_1 x} e^{j\omega t} \hat{a}_y$$

$$\overline{H}_r(x, t) = H_{r0} e^{-\gamma_1 x} e^{j\omega t} \hat{a}_y$$

$$\overline{H}_t(x, t) = H_{t0} e^{-\gamma_2 x} e^{j\omega t} \hat{a}_y$$

Here H_{i0} , H_{r0} and H_{t0} are the magnetic field intensities at $x=0$

Intrinsic impedance for

for perfect dielectric

$$Z = \sqrt{\frac{\mu}{\epsilon}}$$

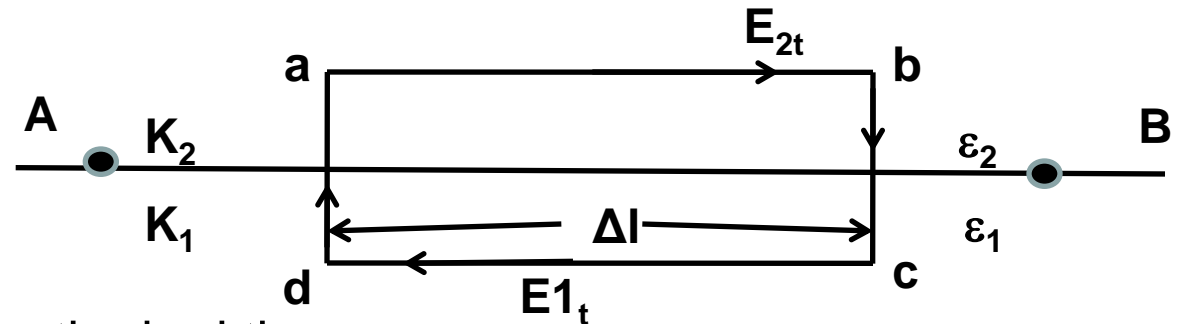
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$$Z = \sqrt{\frac{\mu_0}{\epsilon_0}}$$

Electrostatic Boundary Conditions

Consider a rectangular closed path $abcd$ of infinitesimal area in the plane normal to the boundary and with its side ab and cd parallel to and on either side of the boundary.

Since the line integral $\oint E \cdot dl$ for closed path is zero.



Here considering the closed path $abcd$ then

$$\oint E \cdot dl = \int_a^b E \cdot dl + \int_b^c E \cdot dl + \int_c^d E \cdot dl + \int_d^a E \cdot dl$$

if the loop is shrunk by letting $h \rightarrow 0$

then the parts bc and ad are negligible

$$\oint E \cdot dl \approx \int_a^b E_2 \cdot dl + \int_c^d E_1 \cdot dl = 0$$

E_{1t} and E_{2t} are components of electric field intensities then

$$\oint E \cdot dl = E_{1t} \Delta l - E_{2t} \Delta l = 0$$

$$E_{1t} = E_{2t}$$

Hence any point on the boundary the components of E1 and E2 tangential to the boundary are equal

If medium 1 is conductor then $E_{1t} = 0$ then

$$E_{2t} = 0$$

Boundary Conditions for Displacement vector D

Consider a cylindrical box ABCD which intersects the two dielectric media having permittivities ϵ_1 and ϵ_2 .

Supposing dS is the area of each plane and which is also the area of the boundary enclosed by the cylinder.

Now the charge enclosed by cylinder will be $\rho = \sigma ds$

For finding the boundary conditions for D, apply Gauss's law
Here cylinder will behave like Gaussian surface.

$$\oint D \cdot ds = \sigma ds$$

$$= \int_{upper} D \cdot ds + \int_{lower} D \cdot ds + \int_{curved} D \cdot ds$$

Let $h \rightarrow 0$ then the curved path vanishes

$$\int_{\text{upper}} D \cdot ds + \int_{\text{lower}} D \cdot ds = \int_s \sigma ds$$

